# Compounds with Intermediate Spin. 5. X-ray Study of Tris( $N, N$-dimethyldithiocarbamato)iron(III) at the Two Extreme Temperatures 25 and 400 K 

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#### Abstract

$\mathrm{Fe}\left[\mathrm{S}_{2} \mathrm{CN}\left(\mathrm{CH}_{3}\right)_{2}\right]_{3}$ has a temperature-dependent magnetic moment. Its crystal structure has been determined at 25 and 400 K from X-ray intensities collected with a diffractometer equipped with a Be -walled liquid He cryostat and a microfurnace. The space group is $P b c a$ with $Z=8$ at both temperatures: $a=17 \cdot 288$ (50), $b=$ 19.947 (26), $c=9.998(13) \AA$ at $25 \mathrm{~K}\left(\mu_{\mathrm{eff}} \simeq 2.0\right)$ and $a=17.6961$ (8), $b=20.7906$ (10), $c=10 \cdot 1630$ (6) $\AA$ at $400 \mathrm{~K}\left(\mu_{\text {eff }}=4.83\right)$. The refinements converged to $R=0.046(25 \mathrm{~K})$ and $0.058(400 \mathrm{~K})$. At both temperatures the mononuclear tris(dithiocarbamato) complexes with pseudosymmetry $D_{3}$ are van der Waals packed. The main structural difference between the low- and high-spin complexes is the $\mathrm{Fe}-\mathrm{S}$ distances: the average length is 2.302 (3) at 25 and 2.415 (4) $\AA$ at 400 K . The 25 K low-spin complex is the least distorted. The average apparent decrease in bond lengths between non-H atoms in the planar ligands is only 0.011 (3) $\AA$ when the temperature rises from 25 to 400 K . Preliminary ESCA studies indicate a charge redistribution in the complex: the low-spin $\mathrm{Fe}-\mathrm{S}$ bond is more covalent than the high-spin bond.


## Introduction

Most complexes formed between $\mathrm{Fe}^{111}$ and substituted dithiocarbamate ligands have magnetic moments which are strongly temperature dependent (White, Roper, Kokot, Waterman \& Martin, 1964; Ewald, Martin, Sinn \& White, 1969). The present series of investigations (Albertsson, Elding \& Oskarsson, 1979, and references therein) aims at correlating the magnetic behaviour of various $\mathrm{Fe}\left(\mathrm{S}_{2} \mathrm{CN} R_{2}\right)_{3}$ complexes with the geometrical features of the complexes: (i) the size, symmetry and distortior of the $\mathrm{FeS}_{6}$ core at different temperatures, (ii) bond distances and angles, and conformations of the ligand molecules, especially the $\mathrm{S}_{2} \mathrm{CNC}_{2}$ part, and (iii) the crystal packing of the complexes.

The simplest $\mathrm{Fe}^{\text {III }}$ dithiocarbamate is $\operatorname{tris}(N, N-$ dimethyldithiocarbamato)iron(III). Its crystal structure has previously been investigated at 150 ( $\mu_{\text {eff }}=$ $2.40)^{*}$ and $295 \mathrm{~K}\left(\mu_{\text {eff }}=4.03\right)$ (Albertsson \& Oskarsson, 1977). A Be-walled liquid He cryostat has now been developed for single-crystal diffraction studies down to 20 K on our CAD-4 diffractometer (Albertsson, Oskarsson \& Ståhl, 1979). We have also built a stable microfurnace similar to the one described by Lissalde, Abrahams \& Bernstein (1978) for use in the interval $300-600 \mathrm{~K}$. This equipment has been employed to determine the difference in size of $\mathrm{Fe}^{\mathrm{III}}$ in the low- and high-spin states in a single compound, and how this difference (and the accompanying change in thermal motion of the atoms) affects the distortion and packing of the complex and the geometry of the ligands. Because of crystal decomposition above 400 K ( $\mu_{\text {eff }}=4.83$ ) the pure high-spin state could only be reached by extrapolation. Since high-angle intensities are lacking, the present intensity-data set at 25 K does not warrant a determination of the charge density in the complex, but the result of a preliminary ESCA study is reported.

## Experimental

Single crystals of $\mathrm{Fe}\left[\mathrm{S}_{2} \mathrm{CN}\left(\mathrm{CH}_{3}\right)_{2}\right]_{3}$ were prepared as described by Albertsson \& Oskarsson (1977). Table 1 gives information concerning the crystal data, the collection and reduction of the intensity-data sets and the refinements based on them. The 25 K intensity-data set was collected from three different crystals with the liquid He cryostat. The experimental procedure used is described by Albertsson, Oskarsson \& Ståhl (1979). Since the $\kappa$ motion of the CAD-4 goniostat is limited to $\pm 60^{\circ}$ by the liquid He transfer line and the goniometer head, the three crystals were mounted in different orientations: Nos. 1 and 2 on glass fibres and No. 3 in

[^0]Table 1. Crystal data, collection and reduction of intensities, and the least-squares refinement of $\mathrm{Fe}\left[\mathrm{S}_{2} \mathrm{CN}\left(\mathrm{CH}_{3}\right)_{2}\right]_{3}$
$M_{r}\left(\mathrm{C}_{9} \mathrm{H}_{18} \mathrm{FeN}_{3} \mathrm{~S}_{6}\right)=416.47$, space group: $P b c a, Z=8$. At 25 K (crystals No. $\left.1,2,3\right) a=17.288(50), b=19.947(26), c=9.998(13) \AA$, $V=3348(30) \AA^{3}, D_{x}=1.604 \mathrm{Mg} \mathrm{m}^{-3}$. At 400 K (crystal No. 4) $a=17.6961(8), b=20.7906(10), c=10.1630(6) \AA$, $V=3739$ (6) $\AA^{3}, D_{x}=1.480 \mathrm{Mg} \mathrm{m}^{-3}$.

| Crystal No. | 1 | 2 | 3 | 4 |
| :---: | :---: | :---: | :---: | :---: |
| Crystal size (mm) | $0.06 \times 0.28 \times 0.36$ | $0.12 \times 0.25 \times 0.32$ | $0.06 \times 0.32 \times 0.30$ | $0.12 \times 0.25 \times 0.32$ |
| Radiation | $\mathrm{Cu} K \alpha$ | Cu Ka | Mo $K \alpha$ | $\mathrm{Cu} K \boldsymbol{\alpha}$ |
| Max. $\sin \theta / \lambda\left(\AA^{-1}\right)$ | 0.46 | 0.46 | 1.01 | 0.61 |
| Scan type | $\omega$ | $\omega$ | $\omega$ | $\omega-2 \theta$ |
| Scan width, $\Delta \omega\left(^{\circ}\right.$ ) | $1.0+0.3 \tan \theta$ | $1.0+0.3 \tan \theta$ | $0.8+0.5 \tan \theta$ | $0.7+0.5 \tan \theta$ |
| $\sigma_{c}(I) / I$ requested in a scan | 0.030 | 0.030 | 0.025 | $0 \cdot 020$ |
| Max. recording time (s) | 120 | 120 | 120 | 180 |
| $\mu\left(\mathrm{mm}^{-1}\right)$ | $13 \cdot 52$ | 13.52 | 1.52 | $12 \cdot 51$ |
| Range of transmission factor* | 0.08-0.36 | 0.05-0.23 | 0.74-0.91 | 0.08-0.32 |
| Number of reflexions recorded | 1126 | 615 | 708 | 3979 |
| Number of reflexions in final LS cycles, $m$ | 700 | 380 | 657 | 1295 |
| Number of parameters refined, $n$ | 229 | 229 | 229 | 173 |
| $R=\sum\|\Delta F\| / \sum\left\|F_{o}\right\| \dagger$ | 0.059 | 0.029 | 0.042 | $0 \cdot 058$ |
| overall value |  | 0.046 |  | - |
| $R_{w}=\left[\sum w(\Delta F)^{2} / \sum w\left\|F_{o}\right\|^{2}\right]^{1 / 2}$ | 0.074 | 0.040 | 0.054 | $0 \cdot 065$ |
| overall value |  | 0.061 |  | - |
| $S=\left[\sum w(\Delta F)^{2} /(m-n)\right]^{1 / 2}$ |  | 1.098 |  | 1.070 |
| $C_{1}$ (weighting function) |  | 0.050 |  | $0 \cdot 045$ |
| $C_{2}$ (weighting function) |  | 1.50 |  | - |
| $g \times 10^{-4}$ (extinction) |  | - |  | $0 \cdot 24$ (6) |

an open, thin-walled capillary made of Lindemann glass. For all three crystals the temperature was in the range $22-28 \mathrm{~K}$ during the experiment. Due to mishaps with the electronic temperature controller the temperature variation sometimes increased to more than $\pm 2 \mathrm{~K}$ about the mean, i.e. the usually attainable stability. There are some systematic errors in the structural parameters which are probably caused by the temperature variation, but these errors must be small. At 25 K the zero-point intramolecular movements predominate over all thermal effects in compounds of this type. In order to exclude weak reflexions from the measurements on crystal No. 3, all intensities in the $\sin \theta / \lambda$ range $0.46-1.01 \AA^{-1}$ were calculated with a model based on the data for crystals Nos. 1 and 2.

Crystal No. 4 (Table 1) was mounted in the microfurnace shown in Fig. 1. To improve the stability and allow water cooling to be introduced between the furnace and the $\varphi$-axis shaft of the goniostat, the furnace is mounted on the same type of $x, y, z$ adjustments as is used on the liquid He cryostat. The crystal was secured with 'Superepoxy' cement on a glass fibre. The regulating $\mathrm{Pt} / \mathrm{Pt}(10 \% \mathrm{Rh})$ thermocouple was positioned less than 1 mm below the crystal. The temperature was 400 K , with a long-term stability better than $\pm 1 \mathrm{~K}$.


Fig. 1. Stable, water-cooled microfurnace for CAD-4. The heater section is analogous to that of Lissalde, Abrahams \& Bernstein (1978). The base with $x, y, z$ translation adjustments is of the same type as used with the liquid He cryostat.

During data collection, standard reflexions were checked at regular intervals. No systematic variation in their intensities or in the orientation of the crystals was

Table 2. Positional parameters $\left(\times 10^{5}\right.$, for $\mathrm{H} \times 10^{3}$ ) and approximate isotropic thermal parameters of $\mathrm{Fe}\left[\mathrm{S}_{2} \mathrm{CN}\left(\mathrm{CH}_{3}\right)_{2}\right]_{3}$ at 25 and 400 K , with e.s.d.'s

|  | 25 K |  |  |  | 400 K |  |  |  |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
|  | $x$ | $y$ | $z$ | $B\left(\AA^{2}\right)$ | $x$ | $y$ | $z$ | $B\left(\AA^{2}\right)$ |
| Fe | 56236 (6) | 12194 (5) | 34756 (13) | $0 \cdot 6$ (1) | 56447 (8) | 12217 (7) | 33422 (16) | 4.8 (1) |
| S(1) | 47242 (10) | 7918 (8) | 20197 (21) | 1.1(1) | 46911 (16) | 8295 (14) | 18564 (31) | $6 \cdot 7$ (1) |
| S(2) | 48129 (10) | 21144 (8) | 30354 (21) | $1 \cdot 1$ (1) | 48269 (17) | 21194 (13) | 28218 (36) | 7.3 (1) |
| $\mathrm{N}(1)$ | 37545 (33) | 17616 (30) | 12090 (73) | $1 \cdot 2$ (2) | 37871 (52) | 17739 (58) | 10804 (112) | 7.8 (4) |
| C(1) | 43337 (37) | 15865 (31) | 19637 (84) | 0.9 (2) | 43577 (55) | 16035 (55) | 18113 (113) | $6 \cdot 3$ (4) |
| $\mathrm{C}(2)$ | 33647 (48) | 12975 (39) | 2909 (99) | 1.5 (2) | 33894 (76) | 13299 (89) | 2272 (158) | 11.5 (7) |
| C(3) | 34520 (52) | 24488 (40) | 12203 (100) | 1.6 (2) | 35100 (79) | 24144 (90) | 10275 (172) | 12.7 (7) |
| S(3) | 64018 (9) | 2883 (8) | 32302 (20) | 1.0 (1) | 64273 (15) | 2843 (12) | 30128 (29) | $6 \cdot 0$ (1) |
| S(4) | 65396 (10) | 15395 (8) | 19229 (20) | 1.0 (1) | 66202 (15) | 15243 (12) | 17777 (28) | $5 \cdot 6$ (1) |
| $\mathrm{N}(2)$ | 75900 (31) | 5636 (27) | 15696 (65) | 0.8 (2) | 75944 (46) | 5679 (37) | 14520 (86) | 5.4 (3) |
| C(4) | 69455 (35) | 7580 (31) | 21473 (80) | 0.8 (2) | 69622 (51) | 7540 (45) | 19976 (96) | 4.9 (3) |
| C(5) | 79022 (50) | -1116 (42) | 17797 (113) | 1.6 (2) | 78810 (63) | -807 (57) | 16498 (121) | 7.8 (4) |
| C(6) | 80287 (42) | 9965 (40) | 6706 (101) | 1.2 (2) | 80383 (60) | 9771 (59) | 6127 (123) | 7.5 (4) |
| S(5) | 62930 (10) | 17471 (8) | 51524 (20) | 0.9 (1) | 63151 (14) | 17464 (12) | 50959 (28) | $5 \cdot 5$ (1) |
| S(6) | 50515 (10) | 8360 (8) | 54122 (21) | 1.0 (1) | 50870 (14) | 8478 (13) | 53752 (27) | $6 \cdot 1$ (1) |
| $\mathrm{N}(3)$ | 57781 (35) | 13550 (27) | 75805 (74) | 1.0 (2) | 58161 (52) | 13595 (36) | 73982 (95) | $5 \cdot 3$ (3) |
| C(7) | 57112 (36) | 13251 (31) | 62591 (94) | $1 \cdot 1$ (2) | 57397 (53) | 13282 (42) | 61256 (102) | 4.4 (3) |
| C(8) | 52488 (47) | 10198 (41) | 84833 (90) | 1.2 (2) | 53066 (84) | 10151 (60) | 82875 (120) | 8.9 (5) |
| C(9) | 64038 (44) | 17369 (39) | 82006 (92) | $1 \cdot 1$ (2) | 64088 (81) | 17210 (56) | 80395 (113) | $8 \cdot 0$ (4) |

detected. For crystals Nos. 1 and 2 the intensities were calculated with the peak-location method (Lehmann \& Larsen, 1974) to avoid errors in weak reflexions at low $\theta$ angles. Values of $I$ and $\sigma_{c}(I)$ were corrected for Lorentz, polarization and absorption effects, the latter by numerical integration $\left[\sigma_{c}(I)\right.$ is based on counting statistics].

At both temperatures the cell dimensions were obtained from $\theta$ values determined as $\theta_{h k l}=\left(\omega_{h k l}-\right.$ $\omega_{\overline{\mathrm{h}} \mathrm{k})} / 2$ with the CAD-4 diffractometer in the bisecting mode and measuring $\omega_{i k i}$ at negative $\theta$ angles. The large $\omega$ motions of the standard CAD-4 procedure (DETTH) to measure $\theta$ angles prevent it from being used with the liquid He cryostat or the microfurnace. The cell dimensions at 25 and 400 K are based on 22 and $47 \theta$ values, respectively, measured with Ni -filtered $\mathrm{Cu} K \alpha$ radiation ( $\lambda=1.54056 \AA$ ).

The ESCA experiments were made on an AEI ES200 electron spectrometer by a procedure described elsewhere (Folkesson, 1973).

## Refinements of the $\mathrm{Fe}\left[\mathrm{S}_{2} \mathrm{CN}\left(\mathrm{CH}_{3}\right)_{2}\right]_{3}$ structure at $\mathbf{2 5}$ and 400 K

The coordinates of non-H atoms determined at 295 K were used as initial values in refinement by full-matrix least squares with the structure factors measured at 25 and 400 K . The function minimized was $\sum w(\Delta F)^{2}$ with weights $w=\left(\sigma_{c}^{2} / 4\left|F_{o}\right|^{2}+C_{1}^{2}\left|F_{o}\right|^{2}+C_{2}\right)^{-1} . C_{1}$ and $C_{2}$ were adjusted so that constant values of $\left\langle w(\Delta F)^{2}\right\rangle$ were obtained in different $\left|F_{o}\right|$ and $\sin \theta$ intervals.

Before the last cycles of refinement of the 25 K structure, all its H atoms (belonging to methyl groups)

Table 2 (cont.)
$B$ values for H atoms at 25 K are $1.5 \AA^{2}$, at $400 \mathrm{~K} 8.0 \AA^{2}$.

|  | $x$ | 25 K |  |  | 400 K |  |
| :--- | :---: | ---: | :---: | :---: | ---: | ---: |
|  | $x$ | $y$ | $z$ | $x$ | $z$ |  |
| H(1C2) | $335(5)$ | $148(4)$ | $-56(9)$ | 340 | 142 | -76 |
| H(2C2) | $277(5)$ | $123(4)$ | $53(8)$ | 284 | 123 | 46 |
| H(3C2) | $363(4)$ | $87(4)$ | $39(9)$ | 367 | 88 | 33 |
| H(1C3) | $340(5)$ | $259(4)$ | $216(9)$ | 344 | 260 | 201 |
| H(2C3) | $293(5)$ | $246(4)$ | $81(8)$ | 301 | 249 | 60 |
| H(3C3) | $381(5)$ | $272(4)$ | $93(9)$ | 387 | 275 | 64 |
| H(1C5) | $761(5)$ | $-38(4)$ | $190(9)$ | 751 | -37 | 212 |
| H(2C5) | $821(5)$ | $-28(4)$ | $134(9)$ | 836 | -7 | 227 |
| H(3C5) | $812(5)$ | $-10(4)$ | $267(9)$ | 804 | -27 | 81 |
| H(1C6) | $859(5)$ | $92(4)$ | $63(9)$ | 843 | 125 | 106 |
| H(2C6) | $796(4)$ | $145(4)$ | $91(9)$ | 771 | 129 | 5 |
| H(3C6) | $787(5)$ | $87(4)$ | $-5(9)$ | 824 | 75 | 10 |
| H(1C8) | $501(5)$ | $129(4)$ | $906(9)$ | 541 | 55 | 835 |
| H(2C8) | $490(5)$ | $76(4)$ | $806(9)$ | 476 | 109 | 803 |
| H(3C8) | $558(4)$ | $66(4)$ | $895(8)$ | 537 | 122 | 922 |
| H(1C9) | $623(5)$ | $194(4)$ | $899(10)$ | 618 | 206 | 868 |
| H(2C9) | $658(4)$ | $207(4)$ | $759(9)$ | 669 | 200 | 734 |
| H(3C9) | $679(5)$ | $146(4)$ | $841(9)$ | 666 | 159 | 835 |

were located in a difference map. They were included in the refinement with fixed isotropic temperature factors ( $B=1.5 \AA^{2}$ ). The final positions were then adjusted to fit a difference map of the 400 K structure from geometrical considerations (tetrahedral methyl C atoms with $\mathrm{C}-\mathrm{H}=1.0 \AA$ ). Four of the six methyl groups were rotated about the $\mathrm{C}-\mathrm{N}$ bond. The adjusted positions were included in the structure factor calculations for the 400 K structure (fixed coordinates, $B=8.0 \AA^{2}$ ).

Both data sets were tested for isotropic extinction (Zachariasen, 1967). The correction applied to $\left|F_{o}\right|$ differed from 1.00 for $c a 6 \%$ of the reflexions at 400 K . No extinction was detected in the 25 K set.


Fig. 2. $\delta R$-plot comparison of data and models for $\mathrm{Fe}\left[\mathrm{S}_{2} \mathrm{CN}\left(\mathrm{CH}_{3}\right)_{2}\right]_{3} ; 25 \mathrm{~K}$ : slope 1.012 (4), intercept 0.042 (4); 400 K : slope 0.993 (1), intercept 0.085 (1).

Refinement was terminated when no parameter shift exceeded $5 \%(25 \mathrm{~K}, 10 \%$ for the H parameters) and $1 \%(400 \mathrm{~K})$ of the corresponding e.s.d. Table 2 gives the positional parameters and approximate isotropic temperature factors at both temperatures.* Data and final models are compared by probability plotting in Fig. 2 where ordered values of $\delta R_{i}=\Delta F_{i} / \sigma\left(\left|F_{o}\right|_{i}\right)$ $\left[\sigma\left(\left|F_{0}\right|_{t}\right)=w^{-1 / 2}\right]$ are plotted $v s$ those expected for ordered normal deviates (Abrahams \& Keve, 1971). Both plots are straight lines with intercepts near zero so the systematic errors appear to be small, even for the 25 K data set in spite of the variation in temperature. The slopes are $1.012(4)(25 \mathrm{~K})$ and 0.993 (1) ( 400 K ). This (and the related $S$ values of Table 1) indicates that $\sigma\left(\left|F_{o}\right|\right)$ is on average correctly estimated.

Scattering factors were taken from International Tables for X-ray Crystallography (1974). All factors except those for H were corrected for anomalous dispersion. Since both Mo and $\mathrm{Cu} K \alpha$ radiations were used to measure the structure factors at 25 K , the least-squares program was modified to allow the proper correction of each $\left|F_{c}\right|$.

[^1]
## Description of the $\mathrm{Fe}\left[\mathrm{S}_{2} \mathrm{CN}\left(\mathrm{CH}_{3}\right)_{2}\right]_{3}$ structure at 25 and 400 K

The general features of the structure of $\operatorname{tris}(N, N-$ dimethyldithiocarbamato)iron(III) are the same at 25 and 400 K . Fig. 3 shows stereoscopic pairs of drawings of the mononuclear complex at the two temperatures. Tables 3 and 4 give selected interatomic distances,

Table 3. Interatomic distances ( $\AA$ ) in the $\mathrm{FeS}_{6}$ polyhedron of $\mathrm{Fe}\left[\mathrm{S}_{2} \mathrm{CN}\left(\mathrm{CH}_{3}\right)_{2}\right]_{3}$

|  | 25 K | 400 K |
| :--- | :---: | :---: |
|  | $2.294(4)$ | $2.407(3)$ |
| $\mathrm{Fe}-\mathrm{S}(1)$ | $2.2912(4)$ | $2.420(3)$ |
| $\mathrm{Fe}-\mathrm{S}(2)$ | 2.12. |  |
| $\mathrm{Fe}-\mathrm{S}(3)$ | $2.306(4)$ | $2.414(3)$ |
| $\mathrm{Fe}-\mathrm{S}(4)$ | $2.308(4)$ | $2.430(3)$ |
| $\mathrm{Fe}-\mathrm{S}(5)$ | $2.293(3)$ | $2.403(3)$ |
| $\mathrm{Fe}-\mathrm{S}(6)$ | $2.305(3)$ | $2.418(3)$ |
| $\mathrm{S}(1)-\mathrm{S}(2)$ | $2.831(4)$ | $2.866(4)$ |
| $\mathrm{S}(3)-\mathrm{S}(4)$ | $2.827(4)$ | $2.888(4)$ |
| $\mathrm{S}(5)-\mathrm{S}(6)$ | $2.824(5)$ | $2.880(4)$ |
| $\mathrm{S}(1)-\mathrm{S}(3)$ | $3.299(8)$ | $3.479(4)$ |
| $\mathrm{S}(1)-\mathrm{S}(6)$ | $3.440(5)$ | $3.644(4)$ |
| $\mathrm{S}(3)-\mathrm{S}(6)$ | $3.377(6)$ | $3.572(4)$ |
| $\mathrm{S}(2)-\mathrm{S}(4)$ | $3.386(8)$ | $3.568(4)$ |
| $\mathrm{S}(2)-\mathrm{S}(5)$ | $3.401(6)$ | $3.589(4)$ |
| $\mathrm{S}(4)-\mathrm{S}(5)$ | $3.283(5)$ | $3.446(4)$ |
| $\mathrm{S}(1)-\mathrm{S}(4)$ | $3.76(9)$ | $3.708(4)$ |
| $\mathrm{S}(3)-\mathrm{S}(5)$ | $3.492(4)$ | $3.710(4)$ |
| $\mathrm{S}(6)-\mathrm{S}(2)$ | $3.510(4)$ | $3.733(4)$ |
|  |  |  |



Fig. 3. Stereoscopic drawings of $\mathrm{Fe}\left[\mathrm{S}_{2} \mathrm{CN}\left(\mathrm{CH}_{3}\right)_{2}\right]_{3}$ at 25 and 400 K viewed along the pseudo threefold axis (the complex has approximately $D_{3}$ symmetry). The thermal ellipsoids of the atoms are scaled to include $50 \%$ probability except for the H atoms at 400 K which are drawn as spheres with radii of $0 \cdot 1 \AA$.
angles and torsion angles in the $\mathrm{FeS}_{6}$ core and the ligands. The mean observed $\mathrm{Fe}-\mathrm{S}$ distance increases from 2.302 (3) $\AA$ in the low-spin complex at 25 K to 2.415 (4) $\AA$ in the high-spin complex at 400 K .

There are only van der Waals forces between the complexes. Fig. 4 shows the resulting packing arrangement. There are interactions both between the planar ligands No. 1 and 2 of different complexes, along lines parallel to $x$, and between methyl groups about planes parallel to $2 x-y=0$. There is a slight change in the relative orientation of the complexes when the temperature increases from 25 to 400 K . The closest contacts are across the centres of symmetry with $\mathrm{Fe}-\mathrm{Fe}=$ $6 \cdot 132$ (1) at 25 and 6.511 (1) $\AA$ at 400 K , concurrent with an increase in the molecular volume of the complex from 431 (4) to 467 (1) $\AA^{3}$, which thus reflects both the changing size of the $\mathrm{FeS}_{6}$ core when the spin state changes, but also a more efficient crystal packing at 25 than at 400 K made possible by the very small vibration amplitudes. The magnitudes of the $\beta_{l j}$ values at 25 K are at most only a few times their e.s.d.'s and for $\mathrm{C}(7)$ less than $1 \sigma$. The r.m.s. thermal displacements increase from an average of 0.12 (2) $\AA$ at 25 K to an average of $0.29(5) \AA$ at 400 K . At a slightly higher temperature, about 425 K , the crystals sublime when the intermolecular interactions are balanced by the thermal energy of the complex.

The changes in bond distances and angles in the ligands are very small between the two temperatures. The only appreciable difference is in the S-C-S angles: their mean value increases from 110.6 (1) to

Table 4. Bond distances, bond angles and torsion angles in the three ligands of $\mathrm{Fe}\left[\mathrm{S}_{2} \mathrm{CN}\left(\mathrm{CH}_{3}\right)_{2}\right]_{3}$ ( $\AA$ and deg ) and r.m.s. deviation $(\AA)$ from the leastsquares planes through the $\mathrm{S}_{2} \mathrm{CNC}_{2}$ groups

|  | 25 K | 400 K |
| :---: | :---: | :---: |
| (a) Ligand 1 |  |  |
| $\mathrm{S}(1)-\mathrm{C}(1)$ | 1.724 (7) | 1.715 (12) |
| $\mathrm{S}(2)-\mathrm{C}(1)$ | 1.716 (8) | 1.701 (12) |
| $\mathrm{C}(1)-\mathrm{N}(1)$ | 1.301 (10) | $1 \cdot 303$ (15) |
| $\mathrm{N}(1)-\mathrm{C}(2)$ | 1.468 (11) | 1.449 (20) |
| $\mathrm{N}(1)-\mathrm{C}(3)$ | 1.467 (10) | 1.420 (22) |
| $\mathrm{S}(1)-\mathrm{C}(1)-\mathrm{S}(2)$ | 110.8 (4) | 114.1 (6) |
| $\mathrm{S}(1)-\mathrm{C}(1)-\mathrm{N}(1)$ | 124.5 (6) | 122.5 (9) |
| $\mathrm{S}(2)-\mathrm{C}(1)-\mathrm{N}(1)$ | 124.7 (5) | 123.4 (9) |
| $\mathrm{C}(1)-\mathrm{N}(1)-\mathrm{C}(2)$ | 123.1 (6) | 123.0 (1-2) |
| $\mathrm{C}(1)-\mathrm{N}(1)-\mathrm{C}(3)$ | 121.4 (7) | 123.0 (1.2) |
| $\mathrm{C}(2)-\mathrm{N}(1)-\mathrm{C}(3)$ | 115.5 (7) | 114.0 (1.2) |
| $\mathrm{S}(1)-\mathrm{C}(1) \cdots \mathrm{N}(1)-\mathrm{C}(2)$ | -0.7(1.0) | -0.9(1.7) |
| $\mathrm{S}(1)-\mathrm{C}(1)-\mathrm{N}(1)-\mathrm{C}(3)$ | 179.8 (6) | -179.5 (1.1) |
| $\mathrm{S}(2)-\mathrm{C}(1)-\mathrm{N}(1)-\mathrm{C}(3)$ | -0.3(1.0) | 1.7 (1.8) |
| $\mathrm{S}(2)-\mathrm{C}(1)-\mathrm{N}(1)-\mathrm{C}(2)$ | 179.2 (6) | -179.7(1.0) |
| R.m.s. deviation | 0.0022 | 0.0021 |
| (b) Ligand 2 |  |  |
| S(3)-C(4) | 1.713 (7) | 1.707 (10) |
| S(4)-C(4) | 1.724 (7) | 1.727 (10) |
| $\mathrm{C}(4)-\mathrm{N}(2)$ | 1.313 (9) | $1 \cdot 307$ (12) |
| $\mathrm{N}(2)-\mathrm{C}(5)$ | 1.466 (10) | 1.455 (14) |
| $\mathrm{N}(2)-\mathrm{C}(6)$ | 1.459 (11) | 1.438 (14) |
| $\mathrm{S}(3)-\mathrm{C}(4)-\mathrm{S}(4)$ | 110.7 (4) | 114.5 (5) |
| $\mathrm{S}(3)-\mathrm{C}(4)-\mathrm{N}(2)$ | 125.6 (5) | 124.2 (7) |
| $\mathrm{S}(4)-\mathrm{C}(4)-\mathrm{N}(2)$ | 123.7 (5) | $121 \cdot 3$ (7) |
| $\mathrm{C}(4)-\mathrm{N}(2)-\mathrm{C}(5)$ | 121.4 (6) | 120.9 (9) |
| $\mathrm{C}(4)-\mathrm{N}(2)-\mathrm{C}(6)$ | 122.5 (6) | 123.0 (8) |
| $\mathrm{C}(5)-\mathrm{N}(2)-\mathrm{C}(6)$ | 116.1 (6) | 116.1 (8) |
| $\mathrm{S}(3)-\mathrm{C}(4)-\mathrm{N}(2)-\mathrm{C}(5)$ | 2.0 (1.0) | 2.0 (1.4) |
| $\mathrm{S}(3)-\mathrm{C}(4)-\mathrm{N}(2)-\mathrm{C}(6)$ | -179.1 (6) | -179.2 (8) |
| $\mathrm{S}(4)-\mathrm{C}(4)-\mathrm{N}(2)-\mathrm{C}(6)$ | -0.1(1.0) | -1.6(1.4) |
| $\mathrm{S}(4)-\mathrm{C}(4)-\mathrm{N}(2)-\mathrm{C}(5)$ | -179.0 (6) | 179.6 (8) |
| R.m.s. deviation | 0.0040 | 0.0032 |
| (c) Ligand 3 |  |  |
| S(5)-C(7) | 1.716 (8) | 1.699 (10) |
| S(6)-C(7) | 1.723 (8) | 1.707 (10) |
| $\mathrm{C}(7)-\mathrm{N}(3)$ | 1.328 (12) | $1 \cdot 302$ (14) |
| $\mathrm{N}(3)-\mathrm{C}(8)$ | 1.449 (11) | 1.464 (16) |
| $\mathrm{N}(3)-\mathrm{C}(9)$ | 1.461 (10) | 1.446 (16) |
| S(5)-C(7)-S(6) | 110.4 (5) | 115.5 (6) |
| $\mathrm{S}(5)-\mathrm{C}(7)-\mathrm{N}(3)$ | 124.6 (5) | 121.6 (7) |
| $\mathrm{S}(6)-\mathrm{C}(7)-\mathrm{N}(3)$ | 124.9 (6) | 122.9 (8) |
| $\mathrm{C}(7)-\mathrm{N}(3)-\mathrm{C}(8)$ | 123.0 (6) | 121.7 (9) |
| $\mathrm{C}(7)-\mathrm{N}(3)-\mathrm{C}(9)$ | 120.7 (6) | 123.3 (9) |
| $\mathrm{C}(8)-\mathrm{N}(3)-\mathrm{C}(9)$ | 116.4 (7) | 115.0 (9) |
| $\mathrm{S}(6)-\mathrm{C}(7)-\mathrm{N}(3)-\mathrm{C}(8)$ | -5.8(1.1) | -4.4 (1.4) |
| $\mathrm{S}(6)-\mathrm{C}(7)-\mathrm{N}(3)-\mathrm{C}(9)$ | $174 \cdot 2$ (6) | 174.4 (8) |
| $\mathrm{S}(5)-\mathrm{C}(7)-\mathrm{N}(3)-\mathrm{C}(9)$ | -3.4 (1.0) | -3.1 (1.4) |
| $\mathrm{S}(5)-\mathrm{C}(7)-\mathrm{N}(3)-\mathrm{C}(8)$ | 176.5 (6) | 178.1 (8) |
| R.m.s. deviation | 0.0202 | 0.0160 |

$114.7(4)^{\circ}$. Probability plotting of interatomic distances between the non-H ligand atoms (Albertsson, Elding \& Oskarsson, 1979) indicates that their pooled e.s.d.'s are correctly estimated as could be expected from the $\delta R$ plots of Fig. 2. There is a small but


Fig. 4. Stereoscopic drawings of the crystal packing of $\mathrm{Fe}\left[\mathrm{S}_{2} \mathrm{CN}\left(\mathrm{CH}_{3}\right)_{2}\right]_{3}$ at 25 and 400 K .
significant decrease in the apparent bond lengths when the temperature increases. The average decrease is 0.011 (3) $\AA$ between 25 and 400 K .

The non- H atoms of ligands No. 1 and 2 are planar as shown by both the torsion angles and the r.m.s. deviations from the least-squares planes through the atoms (Table 4). Ligand No. 3 is slightly non-planar at both temperatures. This is most probably caused by the intermolecular packing forces.

## Discussion

Table 5 gives the geometrical characteristics of the $\mathrm{FeS}_{6}$ polyhedron in $\mathrm{Fe}\left[\mathrm{S}_{2} \mathrm{CN}\left(\mathrm{CH}_{3}\right)_{2}\right]_{3}$ at $25,150,295$ and 400 K . The values at 150 and 295 K are taken from Albertsson \& Oskarsson (1977). The mean observed increase in $\mathrm{Fe}-\mathrm{S}$ from its length at 25 K [denoted $\delta(\mathrm{Fe}-\mathrm{S})$ ] is plotted $v s \mu_{\text {eff }}$ in Fig. 5. Linear extrapolation to the spin-only high-spin value of $\mathrm{Fe}^{\mathrm{III}}$, $\mu_{\text {eff }}=5.92$, gives the difference $0.15 \AA$ between the high- and low-spin radii of $\mathrm{Fe}^{3+}$. This difference is much larger than the $0 \cdot 10 \AA$ given by Shannon (1976) in his tabulation of ionic radii, but it is in agreement with the $\mathrm{Fe}-\mathrm{S}$ distance of $2.443 \AA$ in high-spin $\mathrm{Fe}\left(\mathrm{S}_{2} \mathrm{CNC}_{4} \mathrm{H}_{8} \mathrm{O}\right)_{3} \cdot \mathrm{H}_{2} \mathrm{O}$ (Butcher \& Sinn, 1976) and the 2.306 and $2.297 \AA$ in respectively low-spin $\mathrm{Fe}\left[\mathrm{S}_{2} \mathrm{CN}\left(\mathrm{C}_{2} \mathrm{H}_{5}\right)_{2}\right]_{3}$ (Leipoldt \& Coppens, 1973) and low-spin $\mathrm{Fe}\left(\mathrm{S}_{2} \mathrm{CNC}_{4} \mathrm{H}_{4}\right)_{3} \cdot 0 \cdot 5 \mathrm{CH}_{2} \mathrm{Cl}_{2} \quad$ (Bereman, Churchill \& Nalewajek, 1979).

Table 5. The geometry of the coordination polyhedron of $\mathrm{Fe}\left[\mathrm{S}_{2} \mathrm{CN}\left(\mathrm{CH}_{3}\right)_{2}\right]_{3}$ at $25,150,295$ and 400 K

| Temperature (K) | 25 | 150 | 295 | 400 |
| :--- | :--- | :--- | :--- | :--- |
| $\mu_{\text {eff }}$ | - | 2.40 | 4.03 | 4.83 |
| $\mathrm{Fe}-\mathrm{S}(\AA)$ | $2.302(3)$ | $2.339(3)$ | $2.396(5)$ | $2.415(4)$ |
| $\delta(\mathrm{Fe}-\mathrm{S})(\AA)$ | - | $0.036(1)$ | $0.093(3)$ | $0.114(2)$ |
| $\Delta(\AA)$ | 0.191 | 0.218 | 0.248 | 0.269 |
| Ligand bite $(\mathrm{S}-\mathrm{S})(\AA)$ | $2.827(2)$ | $2.843(3)$ | $2.872(4)$ | $2.878(6)$ |
| Edge of triangular face $(\AA)$ | $3.364(25)$ | $3.424(27)$ | $3.517(30)$ | $3.550(30)$ |
| Height of prism $(\AA)$ | $2.474(4)$ | $2.498(3)$ | $2.542(3)$ | $2.554(4)$ |
| Torsion angle $\left({ }^{\circ}\right)$ | $41(1)$ | $40(1)$ | $39(1)$ | $39(2)$ |
| Tilt angle between | $1.6(6)$ | $1.8(8)$ | $1.9(1.0)$ | $2.0(1.0)$ |
| $\quad$ triangular faces $\left({ }^{\circ}\right)$ |  |  |  |  |
| $\mathrm{Fe}-$ centroid of prism $(\AA)$ | $0.027(3)$ | $0.033(2)$ | $0.039(2)$ | $0.036(2)$ |



Fig. 5. The mean observed increase in $\mathrm{Fe}-\mathrm{S}[\delta(\mathrm{Fe}-\mathrm{S})]$ between 25 K and $T$, as a function of $\mu_{\mathrm{eff}}(T)$.

To obtain a single measure of the distortion of the $\mathrm{FeS}_{6}$ polyhedron at the various temperatures, we have calculated the r.m.s. separation $\Delta$ after a best molecular fit of the $\mathrm{FeS}_{6}$ coordinates to those of the ideal octahedron (Dollase, 1974; Albertsson, Elding \& Oskarsson, 1979). At each temperature the mean observed $\mathrm{Fe}-\mathrm{S}$ distance was normalized to $1 \AA$. The least distorted polyhedron is the one at 25 K . When the temperature increases, the $\mathrm{Fe}-\mathrm{S}$ length and the ligand bite ( $\mathrm{S}-\mathrm{S}$ ) increase causing an increase in the height of the polyhedron. The ratio ( $\mathrm{S}-\mathrm{S}$ )/( $\mathrm{Fe}-\mathrm{S}$ ) decreases from 1.228 at 25 to 1.191 at 400 K . Simple electrostatic calculations of ligand-ligand repulsions and the trigonal twist (torsion) made according to Avdeef \& Fackler (1975) predict a decrease in the torsion angle from $41^{\circ}$ at 25 K to $36^{\circ}$ at 400 K . The observed angles have values in this interval but the decrease is not significant.

A ${ }^{13} \mathrm{C}$ NMR study by Gregson \& Doddrell (1975) indicates that the metal-ligand $\pi$-bonding increases when $\mu_{\text {eff }}$ decreases. If there is an accompanying change in the ligand $\mathrm{S}-\mathrm{C}$ and $\mathrm{C}-\mathrm{N}$ lengths it is on a scale below what can be detected from crystal structure determinations by X-ray diffraction methods. We have recorded ESCA electron spectra of the non-H atoms in $\mathrm{Fe}\left[\mathrm{S}_{2} \mathrm{CN}\left(\mathrm{CH}_{3}\right)_{2}\right]_{3}$ at $150,200,250$ and 300 K . The
spectra revealed an increase of the charges on $\mathrm{Fe}, \mathrm{S}$ and N with decreasing temperature: about $0.05-0.1$ a.u. in the interval $300-150 \mathrm{~K}$. There are two types of C atoms in the ligand resulting in broad spectra in which no changes in binding energy could be detected. The ESCA study indicates that in the low-spin complex electron density is moved from the $\mathrm{Fe}, \mathrm{S}$ and N atoms most probably into the $\mathrm{Fe}-\mathrm{S}$ region resulting in a shorter and stronger $\mathrm{Fe}-\mathrm{S}$ bond than in the high-spin complex.

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# Structure du $\mu$-Oxo-bis[dioxo(o-phénanthroline- $N, N^{\prime}$ )-(thiocyanato- $N$ )molybdène(VI)-Acétone 

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#### Abstract

The crystal structure of $\left[\mathrm{Mo}_{2}\left(\mathrm{C}_{12} \mathrm{H}_{8} \mathrm{~N}_{2}\right)_{2}-\right.$ (NCS) $)_{2} \mathrm{O}_{5}$ ]. $\mathrm{C}_{3} \mathrm{H}_{6} \mathrm{O}$ was solved by Patterson and Fourier techniques and refined by full-matrix least-squares methods to a final $R=0.036$ for 2907 independent reflections. The space group is $P 2_{1} / b$ and the unit-cell dimensions are: $a=8.81$ (1), $b=$ 18.64 (2), $c=20.27$ (2) $\AA, \gamma=106.58(8)^{\circ}$, with $Z=4, D_{c}=1.70$ (2), $D_{m}=1.679 \mathrm{Mg} \mathrm{m}^{-3}$. The H

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atoms of the phenanthroline groups were included in structure factor calculations but not refined. Those of the acetone molecule were not searched for. $\left[\mathrm{Mo}_{2}\left(\mathrm{C}_{12} \mathrm{H}_{8} \mathrm{~N}_{2}\right)_{2}(\mathrm{NCS})_{2} \mathrm{O}_{5}\right]$ is a binuclear, oxo-bridged complex of $\mathrm{Mo}^{\mathrm{v1}}$. Each Mo atom has a distorted octahedron of three N and three O atoms. Of the three N atoms, two belong to a phenanthroline and one to a NCS ligand. The $\mathrm{Mo}-\mathrm{O}-\mathrm{Mo}^{\prime}$ bridge is bent. The complex molecule shows a non-negligible torsion around the $\mathrm{Mo}-\mathrm{Mo}^{\prime}$ axis, but the two chelating phenanthroline (C) 1981 International Union of Crystallography


[^0]:    ${ }^{*} \mu_{\text {eff }}$ is the number of Bohr magnetons (1 Bohr magneton $\simeq$ $9.274 \times 10^{-24} \mathrm{~J} \mathrm{~T}^{-1}$ ). © 1981 International Union of Crystallography

[^1]:    * Lists of structure factors and anisotropic thermal parameters (the non-H atoms are assumed to vibrate within the simple harmonic approximation at both temperatures) have been deposited with the British Library Lending Division as Supplementary Publication No. SUP 35596 ( 16 pp.). Copies may be obtained through The Executive Secretary, International Union of Crystallography, 5 Abbey Square, Chester CH1 2HU, England.

